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# Analysis of a mathematical model describing necrotic tumor growth

(joint work with J. Escher and B. Matioc)

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#### Outline



- Motivation
- The mathematical model
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- 3 The moving boundary problem
  - The well-posedness result
  - The transformed problem



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#### Introduction



 We study the growth of a necrotic tumor in different regimes of vascularisation.



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- We study the growth of a necrotic tumor in different regimes of vascularisation.
- The tumor consists of a core of death cells (necrotic core) and a shell of life-proliferating cells surrounding the core (surrounding shell).



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- We study the growth of a necrotic tumor in different regimes of vascularisation.
- The tumor consists of a core of death cells (necrotic core) and a shell of life-proliferating cells surrounding the core (surrounding shell).
- The blood supply provides the nonnecrotic region with nutrients.



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- The concentration of nutrient in the necrotic core is at a constant level which cannot sustain cell proliferation.



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- The necrotic region is not vascularised.



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Conclusions



- We study the growth of a necrotic tumor in different regimes of vascularisation.
- The tumor consists of a core of death cells (necrotic core) and a shell of life-proliferating cells surrounding the core (surrounding shell).
- The blood supply provides the nonnecrotic region with nutrients.
- The concentration of nutrient in the necrotic core is at a constant level which cannot sustain cell proliferation.
- The necrotic region is not vascularised.
- Moreover, no inhibitor chemical species are present



The moving boundary problem







- $\Omega(t)$ -the domain occupied by the nonnecrotic shell
- $\Gamma_1(t)$  outer boundary of the tumor
- $\Gamma_2(t)$  the interior boundary enclosing the necrotic core



The model

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The evolution of the tumor is described by the coupled problem:

•	$\Delta\psi$	=	$\psi$	in $\Omega(t)$ ,	
	Δp	=	0	in $\Omega(t)$ ,	
	$\psi$	=	G	on $\Gamma_1(t)$ ,	
	$\psi$	=	$oldsymbol{G}-\psi_{oldsymbol{0}}$	on $\Gamma_2(t)$ ,	
	p	=	$\kappa_{\Gamma_1(t)} - AG rac{ \mathbf{x} ^2}{4}$	on $\Gamma_1(t)$ ,	(1
	p	=	$\kappa_{\Gamma_2(t)} - AG rac{ \mathbf{x} ^2}{4} - \psi_0$	on $\Gamma_2(t)$ ,	
	$V_i$	=	$\partial_{\nu_i}\psi - \partial_{\nu_i}p - AG \frac{\nu_i \cdot X}{2}$	on $\Gamma_i(t)$ ,	
	Ω(0)	=	Ω <sub>0</sub> ,		

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for  $t \ge 0$  and i = 1, 2.

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- ψ-the rate at which nutrient is added to Ω(t) over the Γ<sub>1</sub>(t) by the vascularization
- p-the pressure inside the tumor





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- p-the pressure inside the tumor
- ν<sub>i</sub>-the restriction of the outward orientated normal at ∂Ω(t) to Γ<sub>i</sub>(t),
- $\kappa_{\Gamma_i}$  the curvature of  $\Gamma_i(t)$
- $V_i$  the normal velocity of  $\Gamma_i$
- x-position vector in R<sup>2</sup>





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- A-describes the balance between the rate of mitosis and apoptosis





- ψ-the rate at which nutrient is added to Ω(t) over the Γ<sub>1</sub>(t) by the vascularization
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- G-the rate of mitosis
- A-describes the balance between the rate of mitosis and apoptosis
- $\psi_0 > 0$ -corresponds to the nutrient concentration assumed constant within the necrotic region
- $\Omega_0$ -the initial tumor domain.



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# Radially symmetric stationary solutions

Theorem (J. Escher, B. Matioc & A. Matioc '10)

Given  $(R_1, R_2) \in (0, \infty)^2$  with  $R_2 < R_1$ , let  $\psi_0^c$  be the constant defined by  $\psi_0^c :=$ 

$$(b_1/R_1 - b_2/R_2) rac{1/R_1 + 1/R_2}{\ln(R_1/R_2)}$$

 $\frac{K_0(R_1)(b_1I_1(R_2)-b_2I_1(R_1))+I_0(R_1)(b_1K_1(R_2)-b_2K_1(R_1))}{I_0(R_1)K_0(R_2)-I_0(R_2)K_0(R_1)}+\frac{R_1^2-R_2^2}{2R_1R_2\ln(R_1/R_2)}$ 

There exists  $A \in \mathbb{R}$  and  $G \in \mathbb{R} \setminus \{0\}$ , such that the annulus

 $A(R_1, R_2) := \{ x \in \mathbb{R}^2 : R_2 < |x| < R_1 \},\$ 

is a stationary solution of problem (1) provided  $\psi_0 \neq \psi_0^c$ . Moreover, A and G are uniquely determined by  $R_1, R_2$ , and  $\psi_0$ . If G = 0, then problem (1) has **no** radially symmetric stationary solutions.



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#### The case G = 0



The annulus A(R<sub>1</sub>, R<sub>2</sub>) centred in zero with radii R<sub>1</sub> > R<sub>2</sub>, is a stationary solution of system (1) iff

$$p'(R_i) = \psi'(R_i), \qquad i = 1, 2,$$

where *p* is the solution of the problem

 $\begin{cases} p'' + \frac{1}{r}p' = 0, & R_2 < r < R_1, \\ p(R_1) = R_1^{-1} - AGR_1^2/4, \\ p(R_2) = -R_2^{-1} - AGR_2^2/4 - \psi_0, \end{cases}$ (2)

when G = 0.



• Given  $G \in \mathbb{R}$ , the solution of (2) is given by the relation  $p(r) = a_{R_1R_2} \ln(r) + b_{R_1R_2}$ ,  $R_2 \le |r| \le R_1$ , with

$$a_{R_1R_2} = \frac{R_1^{-1} + R_2^{-1} + AG(R_2^2 - R_1^2)/4 + \psi_0}{\ln(R_1/R_2)},$$

 $b_{R_1R_2} = R_1^{-1} - AGR_1^2/4 - a_{R_1R_2}\ln(R_1).$ 



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• Given  $G \in \mathbb{R}$ , the solution of (2) is given by the relation  $p(r) = a_{R_1R_2} \ln(r) + b_{R_1R_2}, R_2 \le |r| \le R_1$ , with

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$$b_{R_1R_2} = R_1^{-1} - AGR_1^2/4 - a_{R_1R_2}\ln(R_1).$$

• Furthermore,  $\psi$  is the solution of the problem

$$\begin{cases} \psi'' + \frac{1}{r}\psi' - \psi &= 0, \qquad R_2 < r < R_1, \\ \psi(R_1) &= G, \\ \psi(R_2) &= G - \psi_0, \end{cases}$$
(3)

when G = 0.

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For fixed  $G \in \mathbb{R}$ , the solution of (3) can be written as linear combination of modified Bessel functions of first and second kind

$$\psi = c_{R_1R_2}^1 I_0 + c_{R_1R_2}^2 K_0,$$

with scalars

$$c_{R_1R_2}^1 = \frac{GK_0(R_2) + (\psi_0 - G)K_0(R_1)}{I_0(R_1)K_0(R_2) - I_0(R_2)K_0(R_1)},$$

$$c_{R_1R_2}^2 = \frac{-GI_0(R_2) - (\psi_0 - G)I_0(R_1)}{I_0(R_1)K_0(R_2) - I_0(R_2)K_0(R_1)}.$$



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### NO rad. symm. stat. sol. for G = 0



• Consequently,  $A(R_1, R_2)$  is a steady-state solution of (1) when G = 0 if and only if

$$\frac{\frac{1}{R_1} + \frac{1}{R_2} + \psi_0}{\ln(R_1/R_2)} \frac{1}{R_i} = \psi_0 \frac{K_0(R_1)I_1(R_i) + I_0(R_1)K_1(R_i)}{I_0(R_1)K_0(R_2) - I_0(R_2)K_0(R_1)}, i = 1, 2,$$
(4)

where we used the relations  $I'_0 = I_1$  and  $K'_0 = -K_1$ .



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where we used the relations  $I'_0 = I_1$  and  $K'_0 = -K_1$ .

• It follows then the system (4) has solutions  $(R_1, R_2)$  with  $R_1 > R_2$  exactly when

$$\frac{R_2}{R_1} = \frac{K_0(R_1)I_1(R_1) + I_0(R_1)K_1(R_1)}{K_0(R_1)I_1(R_2) + I_0(R_1)K_1(R_2)}.$$
(5)

• We shown that equality holds in the relation above only when  $R_1 = R_2$  (contradiction).



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#### The case $G \neq 0$

In this case  $A(R_1, R_2)$  is a steady-state solution of (1) exactly when

$$\psi'(R_i) - p'(R_i) - AG\frac{R_i}{2} = 0, \ i = 1, 2.$$
  

$$\Leftrightarrow c_{R_1R_2}^1 I_1(R_i) - c_{R_1R_2}^2 K_1(R_i) - a_{R_1R_2}\frac{1}{R_i} - AG\frac{R_i}{2} = 0, \ i = 1, 2,$$
  

$$\Leftrightarrow a_iG + b_iAG = c_i, \ i = 1, 2, \text{ with}$$



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$$\Leftrightarrow a_i G + b_i AG = c_i, \ i = 1, 2, \text{ with}$$

$$a_{i} := \frac{(K_{0}(R_{2}) - K_{0}(R_{1}))I_{1}(R_{i}) - (I_{0}(R_{1}) - I_{0}(R_{2}))K_{1}(R_{i})}{I_{0}(R_{1})K_{0}(R_{2}) - I_{0}(R_{2})K_{0}(R_{1})},$$

$$b_i := rac{R_1^2 - R_2^2}{4 \ln(R_1/R_2)} rac{1}{R_i} - rac{R_i}{2},$$

 $c_i := -\psi_0 \frac{K_0(R_1)I_1(R_i) + I_0(R_1)K_1(R_i)}{I_0(R_1)K_0(R_2) - I_0(R_2)K_0(R_1)} + \frac{R_1^{-1} + R_2^{-1} + \psi_0}{\ln(R_1/R_2)} \frac{1}{R_i}.$ 

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(6)

#### The case $G \neq 0$

#### Lemma

The system  $a_iG + b_iAG = c_i$ , i = 1, 2, has a (unique) solution (A, G) with  $G \neq 0$  provided that

 $a_1b_2 - a_2b_1 \neq 0,$   $c_1b_2 - c_2b_1 \neq 0,$ and  $c_1 \neq 0$  or  $c_2 \neq 0.$ 



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 The computation done for the case G = 0 shows that c<sub>1</sub> and c<sub>2</sub> cannot be simultaneously zero when R<sub>2</sub> < R<sub>1</sub>.



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- The computation done for the case G = 0 shows that c<sub>1</sub> and c<sub>2</sub> cannot be simultaneously zero when R<sub>2</sub> < R<sub>1</sub>.
- For fixed  $R_1 > 0$  we may see the expression  $a_1b_2 a_2b_1$ as a function of  $R_2 \in (0, R_1)$ . This function is strictly decreasing with respect to  $R_2$ , thus  $a_1b_2 = a_2b_1$  only when  $R_1 = R_2$ .







•  $b_1c_2 = b_2c_1$  if and only if  $\psi_0 = \psi_0^c$ , where

$$\psi_0^c := \frac{(b_1/R_1 - b_2/R_2) \frac{1/R_1 + 1/R_2}{\ln(R_1/R_2)}}{\frac{K_0(R_1)(b_1l_1(R_2) - b_2l_1(R_1)) + l_0(R_1)(b_1K_1(R_2) - b_2K_1(R_1))}{l_0(R_1)K_0(R_2) - l_0(R_2)K_0(R_1)} + \frac{R_1^2 - R_2^2}{2R_1R_2\ln(R_1/R_2)}}{R_1R_2}.$$







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•  $A(R_1, R_2)$  is a stationary solution of (1) if and only if  $\psi \neq \psi_0^c$ and

$$A = \frac{a_1c_2 - a_2c_1}{c_1b_2 - c_2b_1}, \qquad G = \frac{c_1b_2 - c_2b_1}{a_1b_2 - a_2b_1}$$



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## The mathematical settings

 We introduce first a parametrisation for the interfaces Γ<sub>1</sub>(t) and Γ<sub>2</sub>(t),





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#### The mathematical settings

- Let  $0 < R_2 < R_1$  be given and fix  $\alpha \in (0, 1)$ .
- We set  $\mathcal{V} := \{ \rho \in h^{4+\alpha}(\mathbb{S}) : \|\rho\|_{\mathcal{C}(\mathbb{S})} < a \}$ , where

$$a<\frac{R_1-R_2}{R_1+R_2}.$$



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• Each pair  $(\rho_1, \rho_2) \in \mathcal{V}^2$  parametrises a  $C^{4+\alpha}$ -domain  $\Omega(\rho_1, \rho_2) := \left\{ y \in \mathbb{R}^2 : R_2(1 + \rho_2(y/|y|)) < |y| < R_1(1 + \rho_1(y/|y|)) \right\}$ 



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• The condition on *a* ensures that the boundary portions of  $\Omega(\rho_1, \rho_2)$ 

$$\Gamma(\rho_i) := \{ x : |x| = R_i(1 + \rho_i(x/|x|)) \},\$$

i = 1, 2, are disjoint for any choice of  $(\rho_1, \rho_2) \in \mathcal{V}^2$ .



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### The mathematical settings



•  $\Gamma(\rho_i) = N_{\rho_i}^{-1}(0)$ , where  $N_{\rho_i} : \mathbb{R}^2 \setminus \{0\} \to \mathbb{R}, i = 1, 2$ , are defined by

 $N_{
ho_i}(x) = |x| - R_i - R_i 
ho_i (x/|x|), \qquad x \neq 0.$ 



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# The mathematical settings

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eq 0.$$

• The outward unit normal at  $\partial \Omega(\rho_1, \rho_2)$  is given by

$$u_{\rho_1} = \frac{\nabla N_{\rho_1}}{|\nabla N_{\rho_1}|} \text{ on } \Gamma(\rho_1), \text{ and } \nu_{\rho_2} = -\frac{\nabla N_{\rho_2}}{|\nabla N_{\rho_2}|} \text{ on } \Gamma(\rho_2).$$



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 If the function (ρ<sub>1</sub>, ρ<sub>2</sub>) : [0, T] → V<sup>2</sup> describes the motion of the tumor boundaries, then the normal velocity of both boundary components in terms of ρ<sub>i</sub> is given by the formula

$$V_1(t) = -\frac{\partial_t N_{\rho_1}}{|\nabla N_{\rho_1}|}$$
 on  $\Gamma(\rho_1(t))$ , and  $V_2(t) = \frac{\partial_t N_{\rho_2}}{|\nabla N_{\rho_2}|}$  on  $\Gamma(\rho_2(t))$ 

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The sys	ten	n of equations		l l i e 2 L o s 4
$\Delta\psi$	=	$\psi$	in $\Omega(\rho_1, \rho_2)$ ,	$t \ge 0,$
$\Delta p$	=	0	in $\Omega(\rho_1, \rho_2)$ ,	$t \ge 0,$
$\psi$	=	G	on $\Gamma(\rho_1)$ ,	$t \ge 0,$
$\psi$	=	$oldsymbol{G}-\psi_{oldsymbol{0}}$	on $\Gamma(\rho_2)$ ,	$t \ge 0,$
p	=	$\kappa_{\Gamma(\rho_1)} - AG \frac{ x ^2}{4}$	<i>on</i> Γ(ρ <sub>1</sub> ),	$t \ge 0,$
p	=	$\kappa_{\Gamma( ho_2)} - AG rac{ x ^2}{4} - \psi_0$	<b>on</b> Γ(ρ <sub>2</sub> ),	$t \ge 0,$
$\partial_t N_{ ho_i}$	=	$-\langle  abla \psi -  abla p - AGrac{x}{2}    abla$	$\langle N_{ ho_i} \rangle$ on $\Gamma( ho_i)$ ,	<i>t</i> > 0, <i>i</i> = 1, 2,
ρ <sub>1</sub> (0)	=	ρ <sub>01</sub> ,		
ρ <sub>2</sub> (0)	=	ρ02,		(7)

with  $(\rho_1(0), \rho_2(0))$  describing the initial shape of the tumor.

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#### The well-posedness result

#### **Classical solution**

A pair  $(\rho_1, \rho_2, \psi, p)$  is called *classical solution* of (1) on [0, T], T > 0, if

 $\rho_i \in C([0, T], \mathcal{V}) \cap C^1([0, T], h^{1+\alpha}(\mathbb{S})), i = 1, 2,$  $\psi(t, \cdot), \rho(t, \cdot) \in buc^{2+\alpha}(\Omega(\rho_1(t), \rho_2(t))), t \in [0, T],$ 

and if  $(\rho_1, \rho_2, \psi, p)$  solves (7) pointwise.



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and if  $(\rho_1, \rho_2, \psi, p)$  solves (7) pointwise.

#### Theorem (J. Escher, B. Matioc & A. Matioc '10)

Let  $0 < R_2 < R_1$  and  $(A, G, \psi_0) \in \mathbb{R}^3$  be given. There exists an open neighbourhood  $\mathcal{O} \subset \mathcal{V}$  such that for all  $(\rho_1, \rho_2) \in \mathcal{O}^2$ , problem (7) possesses a unique classical solution defined on a maximal time interval  $[0, T(\rho_{01}, \rho_{02}))$  and which satisfies  $(\rho_1, \rho_2)(t) \in \mathcal{O}^2$  for all  $t \in [0, T(\rho_{01}, \rho_{02}))$ .

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#### The problem on the fixed domain



We transform the problem on the fixed domain
 Ω := Ω(0, 0), with boundary Γ<sub>1</sub> := R<sub>1</sub>S and Γ<sub>2</sub> := R<sub>2</sub>S.





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#### The problem on the fixed domain

- Pick  $0 < R_2 < R_1$ ,  $(A, G, \psi_0) \in \mathbb{R}^3$ , and  $\alpha \in (0, 1)$ .
- Given  $(\rho_1, \rho_2) \in \mathcal{V}^2$ , we define  $\Theta_{\rho_1, \rho_2} : \Omega \to \Omega(\rho_1, \rho_2)$  by

$$\Theta_{\rho_1,\rho_2}(x) = \frac{(R_1 - |x|)R_2(1 + \rho_2(x/|x|))}{R_1 - R_2} \frac{x}{|x|} + \frac{(|x| - R_2)R_1(1 + \rho_1(x/|x|))}{R_1 - R_2} \frac{x}{|x|}$$

for  $x \in \Omega$ .

• 
$$\Theta_{\rho_1,\rho_2} \in Diff^{4+\alpha}(\Omega, \Omega(\rho_1, \rho_2))$$

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#### The transformed operators

•  $\mathcal{A}(\rho_1, \rho_2) : buc^{2+\alpha}(\Omega) \to buc^{\alpha}(\Omega)$  $\mathcal{A}(\rho_1, \rho_2)\mathbf{v} := \Delta(\mathbf{v} \circ \Theta_{\rho_1, \rho_2}^{-1}) \circ \Theta_{\rho_1, \rho_2}.$ 



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#### The transformed operators

• 
$$\mathcal{A}(\rho_1, \rho_2) : buc^{2+\alpha}(\Omega) \to buc^{\alpha}(\Omega)$$
  
 $\mathcal{A}(\rho_1, \rho_2)\mathbf{v} := \Delta(\mathbf{v} \circ \Theta_{\rho_1, \rho_2}^{-1}) \circ \Theta_{\rho_1, \rho_2}.$ 

• 
$$\mathcal{B}_i: \mathcal{V}^2 \times (buc^{2+lpha}(\Omega))^2 \to h^{1+lpha}(\mathbb{S})$$
  
 $\mathcal{B}_i(\rho_1, \rho_2, \mathbf{v}, \mathbf{q}) := \frac{1}{R_i} \mathcal{C}_i(\rho_1, \rho_2) \mathbf{v} - \frac{1}{R_i} \mathcal{C}_i(\rho_1, \rho_2) \mathbf{q} - \mathcal{D}_i(\rho_1, \rho_2),$ 



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#### The transformed operators

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$$\mathcal{A}(\rho_1, \rho_2) : buc^{2+\alpha}(\Omega) \to buc^{\alpha}(\Omega)$$
  
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•  $\mathcal{B}_i: \mathcal{V}^2 \times (buc^{2+\alpha}(\Omega))^2 \to h^{1+\alpha}(\mathbb{S})$   $\mathcal{B}_i(\rho_1, \rho_2, \mathbf{v}, \mathbf{q}) := \frac{1}{R_i} \mathcal{C}_i(\rho_1, \rho_2) \mathbf{v} - \frac{1}{R_i} \mathcal{C}_i(\rho_1, \rho_2) \mathbf{q} - \mathcal{D}_i(\rho_1, \rho_2),$ where for  $(\rho_1, \rho_2) \in \mathcal{V}^2$  the linear operators  $\mathcal{C}_i(\rho_1, \rho_2) \in \mathcal{L}(buc^{2+\alpha}(\Omega), h^{1+\alpha}(\mathbb{S})), i = 1, 2, \text{ are given by}$   $\mathcal{C}_i(\rho_1, \rho_2) \mathbf{v}(\mathbf{y}) := \langle \nabla(\mathbf{v} \circ \Theta_{\rho_1, \rho_2}^{-1}) | \nabla \mathbf{N}_{\rho_i} \rangle \circ \Theta_{\rho_1, \rho_2}(R_i \mathbf{y})$ for  $\mathbf{v} \in buc^{2+\alpha}(\Omega)$  and  $\mathbf{y} \in \mathbb{S}.$ 



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#### The transformed operators

• 
$$\mathcal{A}(\rho_1, \rho_2) : buc^{2+\alpha}(\Omega) \to buc^{\alpha}(\Omega)$$
  
 $\mathcal{A}(\rho_1, \rho_2)\mathbf{v} := \Delta(\mathbf{v} \circ \Theta_{\rho_1, \rho_2}^{-1}) \circ \Theta_{\rho_1, \rho_2}.$ 

•  $\mathcal{B}_i: \mathcal{V}^2 \times (buc^{2+\alpha}(\Omega))^2 \to h^{1+\alpha}(\mathbb{S})$   $\mathcal{B}_i(\rho_1, \rho_2, \mathbf{v}, \mathbf{q}) := \frac{1}{R_i} \mathcal{C}_i(\rho_1, \rho_2) \mathbf{v} - \frac{1}{R_i} \mathcal{C}_i(\rho_1, \rho_2) \mathbf{q} - \mathcal{D}_i(\rho_1, \rho_2),$ where for  $(\rho_1, \rho_2) \in \mathcal{V}^2$  the linear operators  $\mathcal{C}_i(\rho_1, \rho_2) \in \mathcal{L}(buc^{2+\alpha}(\Omega), h^{1+\alpha}(\mathbb{S})), i = 1, 2, \text{ are given by}$  $\mathcal{C}_i(\rho_1, \rho_2) \mathbf{v}(\mathbf{y}) := \langle \nabla(\mathbf{v} \circ \Theta_{\rho_1, \rho_2}^{-1}) | \nabla N_{\rho_i} \rangle \circ \Theta_{\rho_1, \rho_2}(R_i \mathbf{y})$ 

for  $v \in buc^{2+\alpha}(\Omega)$  and  $y \in S$ . Moreover,

$$\mathcal{D}_i(
ho_1,
ho_2):=-rac{AG}{R_i}\langlerac{x}{2}|
abla N_{
ho_i}
angle\circ\Theta_{
ho_1,
ho_2}(R_iy).$$



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The model	<b>The radi</b>	ally syl	nmetric case	The moving boundary problem ○○○○○○○●○○○○○○		Conclusions
The <sup>·</sup>	transform	ned	problem		1 1 1 0 2 10 0 4	Leibniz Universität Hannover
ſ	$\mathcal{A}(\rho_1,\rho_2)\mathbf{v}$	=	V		in $\Omega$ ,	$t \ge 0$ ,
	$\mathcal{A}(\rho_1,\rho_2)\boldsymbol{q}$	=	0		<i>in</i> Ω,	$t \ge 0,$
	V	=	G		on $\Gamma_1$ ,	$t \ge 0,$
	V	=	$m{G}-\psi_{0}$		on $\Gamma_2$ ,	$t \ge 0,$
{	q	=	$\frac{1}{R_1}\kappa(\rho_1)-\frac{AGR}{4}$	$\frac{q_1^2}{2}(1+ ho_1)^2$	on $\Gamma_1$ ,	$t \ge 0,$
	q	=	$-\frac{1}{R_2}\kappa(\rho_2)-\frac{A}{2}$	$rac{GR_2^2}{4}(1+ ho_2)^2-\psi_0$	on Γ <sub>2</sub> ,	$t \ge 0,$
	$\partial_t \rho_i$	=	$\mathcal{B}_i(\rho_1,\rho_2,\boldsymbol{v},\boldsymbol{q})$	)	on ₿,	<i>t</i> > 0,
	ρ <sub>1</sub> (0)	=	$ ho_{01},$			
l	ρ <sub>2</sub> (0)	=	$\rho_{02},$			(8)

where  $\mathbf{v} := \psi \circ \Theta_{\rho_1, \rho_2}$ , and  $\mathbf{q} := \mathbf{p} \circ \Theta_{\rho_1, \rho_2}$ ,

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# **Solution operators**

#### Lemma

 $(\rho_1, \rho_2).$ 

Given  $(\rho_1, \rho_2) \in \mathcal{V}^2$ , we let  $\mathcal{T}(\rho_1, \rho_2), \mathcal{S}(\rho_1, \rho_2) \in buc^{2+\alpha}(\Omega)$ denote the unique solution of the Dirichlet problem  $\begin{cases} \mathcal{A}(\rho_1, \rho_2) \mathbf{v} = \mathbf{v} & \text{in } \Omega, \\ \mathbf{v} = \mathbf{G} & \text{on } \Gamma_1, \\ \mathbf{v} = \mathbf{G} - \psi_0 & \text{on } \Gamma_2, \end{cases}$ and  $\begin{cases} \mathbf{A}(\rho_{1},\rho_{2})\mathbf{q} = \mathbf{0} \\ \mathbf{q} = \frac{1}{B_{1}}\kappa(\rho_{1}) - \frac{AGR_{1}^{2}}{4}(1+\rho_{1})^{2} \end{cases}$ in  $\Omega$ , on Γ₁.  $q = -\frac{1}{B_2}\kappa(\rho_2) - \frac{AGR_2^2}{4}(1+\rho_2)^2 - \psi_0$  on  $\Gamma_2$ , respectively. The operators  $\mathcal{T}$  and  $\mathcal{S}$  depend analytically on



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### The evolution equation

• The system (8) reduces to the following evolution equation

 $\partial_t X = \Phi(X) \qquad X(0) = X_0,$  (9)

where  $X := (\rho_1, \rho_2), X_0 := (\rho_{01}, \rho_{02})$ , and  $\Phi := (\Phi_1, \Phi_2)$ .



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# The evolution equation

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$$\partial_t X = \Phi(X) \qquad X(0) = X_0,$$
 (9)

where  $X := (\rho_1, \rho_2), X_0 := (\rho_{01}, \rho_{02})$ , and  $\Phi := (\Phi_1, \Phi_2)$ .

 The components of the nonlocal and nonlinear operator Φ are defined as follows

 $\Phi_i(\rho_1,\rho_2) := \mathcal{B}_i(\rho_i,\mathcal{T}(\rho_1,\rho_2),\mathcal{S}(\rho_1,\rho_2)), \qquad i=1,2.$ 



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# The evolution equation

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 The components of the nonlocal and nonlinear operator Φ are defined as follows

 $\Phi_i(\rho_1,\rho_2) := \mathcal{B}_i(\rho_i,\mathcal{T}(\rho_1,\rho_2),\mathcal{S}(\rho_1,\rho_2)), \qquad i=1,2.$ 

 In order to prove well-posedness of problem (9) it suffices to show that

$$\partial \Phi(0) = \left[ egin{array}{cc} \partial_{
ho_1} \Phi_1(0) & \partial_{
ho_2} \Phi_1(0) \ \partial_{
ho_1} \Phi_2(0) & \partial_{
ho_2} \Phi_2(0) \end{array} 
ight]$$

generates a strongly continuous and analytic semigroup.



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#### Theorem

The operator  $\Phi$  is analytic, i.e.  $\Phi \in C^{\omega}(\mathcal{V}^2, (h^{1+\alpha}(\mathbb{S}))^2)$ . The Fréchet derivative  $\partial \Phi(0)$ , seen as an unbounded operator in  $(h^{1+\alpha}(\mathbb{S}))^2$  with domain  $(h^{4+\alpha}(\mathbb{S}))^2$  generates a strongly continuous and analytic semigroup in  $\mathcal{L}((h^{1+\alpha}(\mathbb{S}))^2)$ , i.e.

 $-\partial \Phi(\mathbf{0}) \in \mathcal{H}((h^{4+\alpha}(\mathbb{S}))^2, (h^{1+\alpha}(\mathbb{S}))^2).$ 



#### Theorem

The operator  $\Phi$  is analytic, i.e.  $\Phi \in C^{\omega}(\mathcal{V}^2, (h^{1+\alpha}(\mathbb{S}))^2)$ . The Fréchet derivative  $\partial \Phi(0)$ , seen as an unbounded operator in  $(h^{1+\alpha}(\mathbb{S}))^2$  with domain  $(h^{4+\alpha}(\mathbb{S}))^2$  generates a strongly continuous and analytic semigroup in  $\mathcal{L}((h^{1+\alpha}(\mathbb{S}))^2)$ , i.e.

 $-\partial \Phi(\mathbf{0}) \in \mathcal{H}((h^{4+\alpha}(\mathbb{S}))^2, (h^{1+\alpha}(\mathbb{S}))^2).$ 

#### Proof

• 
$$\partial_{\rho_1} \Phi_1(0)[\rho_1] = A_{11} + B_{11}$$
, where  
 $B_{11} \in \mathcal{L}(h^{2+\alpha}(\mathbb{S}), h^{1+\alpha}(\mathbb{S}))$   
 $A_{11\rho_1} := \frac{1}{R_1^2} C_1(0)(\Delta, \operatorname{tr}_1, \operatorname{tr}_2)^{-1}(0, \rho_1'', 0), \quad \forall \rho_1 \in h^{4+\alpha}(\mathbb{S})$   
 $A_{11\rho_1}(y) = -\frac{1}{R_1^3} \sum_{m \in \mathbb{Z} \setminus \{0\}} \frac{R_1^{|m|} R_2^{-|m|} + R_1^{-|m|} R_2^{|m|}}{R_1^{|m|} R_2^{-|m|} - R_1^{-|m|} R_2^{|m|}} |m|^3 \widehat{\rho}_1(m) y^m,$   
for  $\rho_1(y) = \sum_m \widehat{\rho}_1(m) y^m.$ 

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• 
$$\partial_{\rho_2} \Phi_1(0) = A_{12} + B_{12}$$
, where  
 $B_{12} \in \mathcal{L}(h^{2+\alpha}(\mathbb{S}), h^{1+\alpha}(\mathbb{S}))$   
 $A_{12}\rho_2 := -\frac{1}{R_1R_2}C_1(0)(\Delta, \operatorname{tr}_1, \operatorname{tr}_2)^{-1}(0, 0, \rho_2') \quad \forall \rho_2 \in h^{4+\alpha}(\mathbb{S})$   
 $A_{12}\rho_2(y) = -\frac{1}{R_1^2R_2}\sum_{m \in \mathbb{Z} \setminus \{0\}} \frac{2}{R_1^{|m|}R_2^{-|m|} - R_1^{-|m|}R_2^{|m|}} |m|^3 \widehat{\rho}_2(m) y^m$ ,  
rovided that  $\rho_2 = \sum_{m \in \mathbb{Z}} \widehat{\rho}_2(m) y^m$ .



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$$\partial_{\rho_2} \Phi_1(0) = A_{12} + B_{12}$$
, where  
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 $A_{12\rho_2} := -\frac{1}{R_1R_2}C_1(0)(\Delta, \operatorname{tr}_1, \operatorname{tr}_2)^{-1}(0, 0, \rho_2'') \quad \forall \rho_2 \in h^{4+\alpha}(\mathbb{S})$   
 $A_{12\rho_2}(y) = -\frac{1}{R_1^2R_2}\sum_{m \in \mathbb{Z} \setminus \{0\}} \frac{2}{R_1^{|m|}R_2^{-|m|} - R_1^{-|m|}R_2^{|m|}} |m|^3 \widehat{\rho}_2(m) y^m$ ,  
rovided that  $\rho_2 = \sum_{m \in \mathbb{Z}} \widehat{\rho}_2(m) y^m$ .  
•  $\partial_{\rho_2} \Phi_2(0) = A_{22} + B_{22}$ , with  
 $B_{22} \in \mathcal{L}(h^{2+\alpha}(\mathbb{S}), h^{1+\alpha}(\mathbb{S}))$   
 $A_{22\rho_2} := -\frac{1}{R_2^2} C_2(0)(\Delta, \operatorname{tr}_1, \operatorname{tr}_2)^{-1}(0, 0, \rho_2'') \quad \forall \rho_2 \in h^{4+\alpha}(\mathbb{S})$   
 $A_{22\rho_2}(y) = -\frac{1}{R_2^3} \sum_{m \in \mathbb{Z} \setminus \{0\}} \frac{R_1^{|m|}R_2^{-|m|} + R_1^{-|m|}R_2^{|m|}}{R_1^{|m|}R_2^{-|m|} - R_1^{-|m|}R_2^{|m|}} |m|^3 \widehat{\rho}_2(m) y^m$ 

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• 
$$\partial_{\rho_1} \Phi_2(0) = A_{21} + B_{21}$$
, where

$$\begin{split} B_{21} &\in \mathcal{L}(h^{2+\alpha}(\mathbb{S}), h^{1+\alpha}(\mathbb{S})) \\ A_{21}\rho_1 &:= \frac{1}{R_1R_2}C_2(0)(\Delta, \mathrm{tr}_1, \mathrm{tr}_2)^{-1}(0, \rho_1'', 0) \quad \forall \rho_2 \in h^{4+\alpha}(\mathbb{S}) \\ A_{21}\rho_1(y) &= -\frac{1}{R_1R_2^2}\sum_{m \in \mathbb{Z}\setminus\{0\}} \frac{2}{R_1^{|m|}R_2^{-|m|} - R_1^{-|m|}R_2^{|m|}} |m|^3 \widehat{\rho}_1(m) y^m, \end{split}$$

for all functions  $\rho_1 = \sum_{m \in \mathbb{Z}} \widehat{\rho}_1(m) y^m$  in  $h^{4+\alpha}(\mathbb{S})$ .



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• 
$$\partial_{\rho_1} \Phi_2(0) = A_{21} + B_{21}$$
, where

$$\begin{split} B_{21} &\in \mathcal{L}(h^{2+\alpha}(\mathbb{S}), h^{1+\alpha}(\mathbb{S})) \\ A_{21}\rho_1 &:= \frac{1}{R_1R_2}C_2(0)(\Delta, \mathrm{tr}_1, \mathrm{tr}_2)^{-1}(0, \rho_1'', 0) \quad \forall \rho_2 \in h^{4+\alpha}(\mathbb{S}) \\ A_{21}\rho_1(y) &= -\frac{1}{R_1R_2^2}\sum_{m \in \mathbb{Z} \setminus \{0\}} \frac{2}{R_1^{|m|}R_2^{-|m|} - R_1^{-|m|}R_2^{|m|}} |m|^3 \widehat{\rho}_1(m) y^m, \end{split}$$

for all functions  $\rho_1 = \sum_{m \in \mathbb{Z}} \widehat{\rho}_1(m) y^m$  in  $h^{4+\alpha}(\mathbb{S})$ .

 The operators A<sub>ij</sub>, 1 ≤ i, j ≤ 2, found above are all Fourier multipliers, since they are of the form

$$\sum_{m\in\mathbb{Z}}\widehat{\rho}(m)y^m\mapsto\sum_{m\in\mathbb{Z}}M_k\widehat{\rho}(m)y^m$$

with symbol  $(M_k)_{k \in \mathbb{Z}} \subset \mathbb{C}$ .



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### The matrix $\partial \Phi(0)$ is a generator

•  $-A_{ii} \in \mathcal{H}(h^{4+\alpha}(\mathbb{S}), h^{1+\alpha}(\mathbb{S})), i = 1, 2.$ 



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#### The matrix $\partial \Phi(0)$ is a generator

• 
$$-A_{ii} \in \mathcal{H}(h^{4+\alpha}(\mathbb{S}), h^{1+\alpha}(\mathbb{S})), i = 1, 2.$$

•  $A_{12}, A_{21} \in \mathcal{L}(h^{2+\alpha}(\mathbb{S})).$ 



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#### The matrix $\partial \Phi(0)$ is a generator



- $-A_{ii} \in \mathcal{H}(h^{4+\alpha}(\mathbb{S}), h^{1+\alpha}(\mathbb{S})), i = 1, 2.$
- $A_{12}, A_{21} \in \mathcal{L}(h^{2+\alpha}(\mathbb{S})).$
- $h^{2+\alpha}(\mathbb{S}) = (h^{1+\alpha}(\mathbb{S}), h^{4+\alpha}(\mathbb{S}))_{1/3}.$



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#### The matrix $\partial \Phi(0)$ is a generator



- $-A_{ii} \in \mathcal{H}(h^{4+\alpha}(\mathbb{S}), h^{1+\alpha}(\mathbb{S})), i = 1, 2.$
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- $h^{2+\alpha}(\mathbb{S}) = (h^{1+\alpha}(\mathbb{S}), h^{4+\alpha}(\mathbb{S}))_{1/3}.$
- $-\partial_{\rho_i}\Phi_i(0) \in \mathcal{H}(h^{4+\alpha}(\mathbb{S}), h^{1+\alpha}(\mathbb{S})), i = 1, 2,$

while the elements on the secondary diagonal belong to  $\mathcal{L}(h^{2+\alpha}(\mathbb{S}), h^{1+\alpha}(\mathbb{S}))$ , and having thus lower order.



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### The matrix $\partial \Phi(0)$ is a generator



- $-A_{ii} \in \mathcal{H}(h^{4+\alpha}(\mathbb{S}), h^{1+\alpha}(\mathbb{S})), i = 1, 2.$
- $A_{12}, A_{21} \in \mathcal{L}(h^{2+\alpha}(\mathbb{S})).$
- $h^{2+\alpha}(\mathbb{S}) = (h^{1+\alpha}(\mathbb{S}), h^{4+\alpha}(\mathbb{S}))_{1/3}.$
- $-\partial_{\rho_i}\Phi_i(\mathbf{0})\in\mathcal{H}(h^{4+lpha}(\mathbb{S}),h^{1+lpha}(\mathbb{S})),i=1,2,$

while the elements on the secondary diagonal belong to  $\mathcal{L}(h^{2+\alpha}(\mathbb{S}), h^{1+\alpha}(\mathbb{S}))$ , and having thus lower order.

Thus the matrix

$$\partial \Phi(\mathbf{0}) = \left[ egin{array}{cc} \partial_{
ho_1} \Phi_1(\mathbf{0}) & \partial_{
ho_2} \Phi_1(\mathbf{0}) \ \partial_{
ho_1} \Phi_2(\mathbf{0}) & \partial_{
ho_2} \Phi_2(\mathbf{0}) \end{array} 
ight]$$

generates a strongly continuous and analytic semigroup, which completes the proof.



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- - We study a model describing the growth of necrotic tumors in different regimes of vascularisation.





- We study a model describing the growth of necrotic tumors in different regimes of vascularisation.
- We determine all radially symmetric stationary solutions and reduce the moving boundary problem into a nonlinear evolution equation for the functions parameterising the boundaries of the shell.





- We study a model describing the growth of necrotic tumors in different regimes of vascularisation.
- We determine all radially symmetric stationary solutions and reduce the moving boundary problem into a nonlinear evolution equation for the functions parameterising the boundaries of the shell.
- Parabolic theory provides a suitable context for proving local well-posedness of the problem for small initial data.



#### Conclusions

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#### Thank you for your attention!

